

# Modular Invariants and Fischer-Griess Monster <sup>1</sup>

**Marcin Jankiewicz<sup>†</sup>, Thomas W. Kephart<sup>†</sup>**

<sup>†</sup> Department of Physics and Astronomy, Vanderbilt University, Nashville, TN 37235, USA

**Abstract.** We show interesting relations between extremal partition functions of a family of conformal field theories and dimensions of the irreducible representations of the Fischer-Griess Monster sporadic group. We argue that these relations can be interpreted as an extension of Monster moonshine.

## 1 Introduction

Partition functions of conformal field theories with central charge  $c = 24 \cdot k$ , where  $k = 1, 2, \dots$ , are modular invariants. There exists a relation between CFT with  $c = 24$ , i.e.  $k = 1$ , and the Monster sporadic group. This is the famous Monster moonshine theorem that states that coefficients of the  $q$ -expansion are related in a simple manner to the dimensions of the irreducible representations of the this largest sporadic group. Here, we are going conjecture a natural extension of this relation to the higher dimensional, i.e.  $k > 1$ , case. We are going to use lattices, or more precisely their  $\Theta$ -functions to describe a given CFT. The  $q$ -expansion of a  $\Theta$ -function of a lattice  $\Lambda$  is given as

$$Z_\Lambda = \sum_{x \in \Lambda} N(m) q^m, \quad (1)$$

where we sum over all vectors  $x$ , in the lattice  $\Lambda$ , with length  $m = x \cdot x$ .  $N(m)$  is the number of vectors of norm  $m$  and  $q \equiv e^{i\pi\tau}$  where  $\tau$  is the modular parameter. The spectra of meromorphic conformal field theories can be expressed in terms of partition functions of even self-dual lattices. This means that the exponent  $m$  in the  $q$ -expansion will be necessarily an even number. Both self-duality and evenness of a lattice correspond to invariance of a partition function  $\mathcal{Z}$  (which is closely related to  $Z_\Lambda$ ) under the generators  $S$  and  $T$  of a modular group  $SL(2, \mathbb{Z})$ .

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Formally,  $Z_\Lambda$  of a  $d$  dimensional lattice  $\Lambda$  is a modular form of weight  $d/2$ . The partition function  $\mathcal{Z}$  of a lattice is defined as follows

$$\mathcal{Z} = Z_\Lambda / \eta^{d/2}, \quad (2)$$

where  $\eta(q) = q^{1/12} \prod_{m=1}^{\infty} (1 - q^{2m})$  is Dedekind  $\eta$ -function which is a modular form of weight  $1/2$ . The partition functions of all the 24 dimensional even self-dual lattices (the Niemeier lattices) can be written as

$$\mathcal{Z} = [J + 24(h+1)] \eta^{24}, \quad (3)$$

where  $h$  is the Coxeter number of a given lattice. For example,  $h = 0$  corresponds to the famous Leech lattice,  $h = 30$  to the Niemeier lattice based on a root system of  $E_8^3$ , etc. Physically  $24(h+1)$  corresponds to a number of massless states in a given theory. Using a technique presented in [1, 2], one can choose any Niemeier lattice  $\Lambda_1$  to generate the  $\Theta$ -function of another Niemeier lattice  $\Lambda_2$ . In [1] it was shown that it is possible to generate extremal partition functions [3] by taking the  $k^{\text{th}}$  power of (3) and treating  $x_i = 24(h_i + 1)$  (where  $i = 1, \dots, k$ ) as a free parameter.

## 2 CFTs with $c = 24 \cdot k$ : Systematic Approach

Different choices of constant parameters  $x_i$ s correspond to different  $\Theta$ -functions. One can find corresponding extremal partition functions, (2) and write them as  $q$ -expansions of the form

$$\prod_{i=1}^k (J + 24 + x_i) = \frac{1}{q^{2k}} \left[ 1 + \sum_{m=(k-1)}^{\infty} f_{2m}(x_1, \dots, x_k) q^{2m-2k} \right] \quad (4)$$

As an example, we will work with a choice of  $k-1$  parameters  $x_i$ s that eliminates coefficients of all but one term with negative powers in the  $q$ -expansion that correspond (in a field theoretic language) to tachyonic states. In this setup we are left with only one free parameter  $x_k$ . Different choices of  $x_k$  would correspond to different partition function candidates for conformal field theories with  $c = 24 \cdot k$ . Here we list the first three cases:

$$\mathcal{G}_1(x_1) = \frac{1}{q^2} + (24 + x_1) + 196884q^2 + \dots \quad (5a)$$

$$\mathcal{G}_2(x_2) = \frac{1}{q^4} + (393192 - 48x_2 - x_2^2) + 42987520q^2 + \dots \quad (5b)$$

$$\mathcal{G}_3(x_3) = \frac{1}{q^6} + (50319456 - 588924x_3 + 72x_3^2 + x_3^3) + 2592899910q^2 + \dots \quad (5c)$$

Notice that in each case all of the tachyonic states (except the lowest one) are absent. Since the allowed values [4] of the coefficient of  $q^0$  are integers that run from zero to the value of the  $q^2$  coefficient, one can easily find the number of “allowed” partition functions in  $24k$  dimensions [1].

### 3 Monster Moonshine and its Extension

The extremal 24 dimensional case has been shown to be related to the Fischer-Griess monster group. In mathematics this fact is known as Monster moonshine ([5] and [6]). One can evaluate  $\mathcal{G}_1$  at  $x_1 = -24$  which corresponds to the  $j$ -invariant to find

$$j = \frac{1}{q^2} + 196884q^2 + 21493760q^4 + 864299970q^6 + 20245856256q^8 + \dots \quad (6)$$

The coefficients of this expansion decompose into dimensions of the irreducible representations of the Monster<sup>2</sup>, where we use the notation  $j = \frac{1}{q^2} + j_2q^2 + j_4q^4 + \dots$ . Following this interpretation of the Monster, one can easily generate moonshine in higher dimensional cases, i.e., one can express coefficients of any partition functions, for example  $\mathcal{G}_k(x_k)$ , for any choice of  $k$ , in terms of the dimensions of irreducible representations of the Monster group. We present a few of our results [1] in Table-1, where coefficients of  $G_k(x_k)$  are expressed in terms of the coefficients of the invariant function  $j$ , that (via the original Monster moonshine) are related to the Monster. We notice [1] that the coefficients  $g_{2n}$  fall into patterns with period  $k!$ , and conjecture that this periodicity also continues to hold for all  $k$ . The polynomial conditions to be satisfied to find the extremal partition functions for large  $k$  become increasingly more difficult to solve with increasing  $k$ , so we do not have results for  $k > 6$ .

Table-1 give the general periodicity in coefficients of  $\mathcal{G}_k(x_k)$  for  $k < 6$ . To summarize, when  $k = 1$  it is known via standard Monster Moonshine that the coefficients of  $j$  decompose into Monster representations [6]. The fact that all the higher  $k$  coefficients also decompose into Monster representations indicates that they have large symmetries containing the Monster and the fact that they have these symmetries may indicate that they are related to  $24k$  dimensional lattices.

<sup>2</sup> for explicit realization of the Monster moonshine see [1].

$k = 2$	$k = 2$	$k = 3$	$k = 3$
$g_{4i+2}$	$2j_{2(4i+2)}$	$g_{6i+2}$	$3j_{3(6i+2)}$
$g_{4i+4}$	$2j_{2(4i+4)} + j_{2(2i+2)}$	$g_{6i+4}$	$3j_{3(6i+4)}$
		$g_{6i+6}$	$3j_{3(6i+6)} + j_{2i+2}$
$k = 4$	$k = 4$	$k = 5$	$k = 5$
$g_{8i+2}$	$4j_{4(8i+2)}$	$g_{10i+2}$	$5j_{5(10i+2)}$
$g_{8i+4}$	$4j_{4(8i+4)} + 2j_{2(2i+4)}$	$g_{10i+4}$	$5j_{5(10i+4)}$
$g_{8i+6}$	$4j_{4(8i+6)}$	$g_{10i+6}$	$5j_{5(10i+6)}$
$g_{8i+8}$	$4j_{4(8i+8)} + 2j_{(8i+8)} + j_{2i+2}$	$g_{10i+8}$	$5j_{5(10i+8)}$
		$g_{10i+10}$	$5j_{5(10i+10)} + j_{2i+2}$

**Table 1.** Periodicity of the coefficients  $g_n$  for  $c = 24 \cdot k$  extremal partition functions  $\mathcal{G}_k$  in terms of coefficients  $j_{2n}$  of the modular function  $j$ .

#### 4 Final Comments

We have shown that there exists a family of transformations, more general than of simple  $\mathbb{Z}_2$  form, that relates the members of the class of holomorphic conformal field theories, i.e., the Niemeier lattices. Furthermore these transformations connect non-Niemeier modular invariant  $c = 24$   $\Theta$ -functions. These results generalize to any  $c = 24 \cdot k$  case (in particular, when the resulting parametrization corresponds to  $24k$  dimensional lattice). In 24 dimensions there is one extremal partition function, corresponding to a Leech lattice, whose  $q$ -expansion coefficients can be written in the form of a linear combinations of dimensions of the irreducible representations of the Monster group; this relationship is at the core of the Monster Moonshine. The decomposition is possible only in the extremal case, since in all other cases a non-zero constant term in the  $q$ -expansion would be present. The presence of this term in the expansion would imply the introduction of an enormous unnatural set of singlets in the decomposition. We believe a similar situation occurs at  $24k$  where we have extended this argument. Instead of a Leech lattice we have to deal with higher dimensional extremal partition functions (but note [1] that there is more than one type of extremal partition function in  $24k$  dimensions). Existence of  $24k$  extremal lattices for  $k > 2$  is only a conjecture [3], but our results are consistent with and provides supporting evidence for this conjecture. One can use our generalized version of Monster Moonshine to postulate that we already have the  $q$ -expansion of higher dimensional extremal lattices, and that the symmetry provided by the Monster decomposition can be used to learn more about these lattices.

We suggest that the extremal  $\mathcal{G}$  partition functions will generate new  $c = 24 \cdot k$  CFTs. We know the first at  $k = 1$  corresponds to the Leech lattice, the second  $k = 2$  case also corresponds to the known lattice,  $P_{48}$ , and since the higher  $k$  cases all possess Monster symmetry we conjecture that it is likely that they correspond to CFTs constructed on extremal lattices in  $24k$ -dimensions. To the best of our knowledge, Monster symmetry was not known to come into play except at  $k = 1$ .

Using the techniques presented in [1], one can construct a large class of conformal field theories with central charge that is a multiple of 24. We have demonstrated (or at least conjectured) the possibility of a new realization of Monster moonshine. This is realized as a periodicity in a pattern of coefficients in  $q$ -expansions of the extremal partition functions.

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